

Lemma 2 Falls $\lambda > \mu$, dann gilt $\forall x \in \mathbb{C}S_n$:

$$p_\lambda \times q_\mu = 0 = q_\mu \times p_\lambda$$

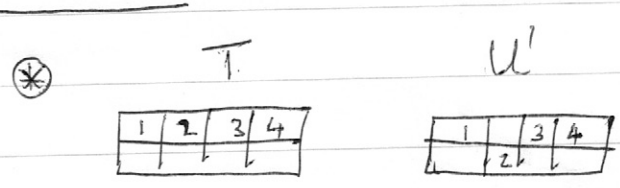
Inbesondere gilt $c_\lambda c_\mu = 0$

(Es gilt auch $c_\mu c_\lambda = 0$ ~~nicht~~ ~~Feld~~ ~~Homis~~, also das brauchen wir nicht.)

Bew: Wir zeigen $p_\lambda \sigma q_\mu = 0 \quad \forall \sigma \in S_n$. Sei dazu T wie gehabt das ^{kanonische} ~~Standard~~ ~~Tableau~~ des Form λ und U das ^{kanonische} ~~Standard~~ ~~Tableau~~ des Form μ . Dann ist $\sigma U = U'$ ein weiteres Tableau des Form μ und $\sigma q_\mu \sigma^{-1} = q_{\mu, \sigma}$. Es genügt also zu zeigen, dass $p_\lambda q_{\mu, \sigma} = 0$ für jedes Tableau $\sigma U = U'$ des Form μ .

Tauschprinzip 2 Falls $\lambda > \mu$, so existieren Elemente $a, b \in \mathbb{N}$ die in der gleichen Zeile von T und in der gleichen Spalte von U' vorkommen (*). Die Transposition τ von a und b erfüllt dann

$$p_\lambda q_{\mu, \sigma} \stackrel{\tau}{=} p_\lambda q_{\mu, \sigma} \stackrel{\tau}{=} p_\lambda \tau q_{\mu, \sigma} = -p_\lambda q_{\mu, \sigma} = 0$$



Herausstreichen der langen Zeilen von T aus U' geht solange gut, wie $\lambda_i = \mu_i$. $\lambda_i > \mu_i$ erzwingt dann doppelt besetzte Spalten in U' .

We are now ready to prove the main theorem.

• First, we claim that $c_\lambda \cdot c_\lambda = W c_\lambda$.

Indeed, this follows from the fact that $\forall \pi \in P, \exists \xi \in Q$

$$\pi c_\lambda c_\lambda \xi = \rho(\pi(\xi)) c_\lambda c_\lambda$$

so Lemma 1 implies $\exists W$ s.t. $c_\lambda^2 = W c_\lambda$.

To show that $W \neq 0$, let's calculate it. Note that the trace of right multiplication by c_λ on $\mathbb{C}S_n$ is $n!$ (since $c_\lambda(e) = 1$).

On the other hand c_λ maps any $x c_\lambda \in \mathbb{C}S_n c_\lambda$ to W times itself and anything in the complement into $\mathbb{C}S_n c_\lambda$. So the trace on $\mathbb{C}S_n = \mathbb{C}S_n c_\lambda + \text{complement}$ is $\dim V_\lambda \cdot W$. So

$$W = \frac{n!}{\dim V_\lambda}$$

• To show that V_λ are irreducible, ~~we~~ assume $W \subset V_\lambda$ is ~~invariant~~ invariant subspace, and consider $c_\lambda W$. Since $c_\lambda V_\lambda = \mathbb{C} c_\lambda$ by Lemma 4, we must have either $c_\lambda W = \mathbb{C} c_\lambda$ or $c_\lambda W = 0$.

If $c_\lambda W = \mathbb{C}c_\lambda$, $\exists w \in W$ st. $c_\lambda w = c_\lambda$.

But since any $v \in V_\lambda$ is of the form $v = xc_\lambda$
we learn

$$v = xc_\lambda = xc_\lambda w \in W.$$

So $V_\lambda = W$.

If $c_\lambda W = 0$, then $c_\lambda W^\perp = \mathbb{C}c_\lambda$, where W^\perp is complementary invariant subspace. But then by the argument we just gave $W^\perp = V$, so $W = 0$.

• Finally, we show that if $\lambda > \mu$, then V_λ and V_μ are not isomorphic. Indeed, $c_\lambda V_\lambda = \mathbb{C}c_\lambda$ but $c_\lambda V_\mu = p_\lambda(q_\lambda \mathbb{C} \begin{smallmatrix} \sim \\ \sim \end{smallmatrix} p_\mu)q_\mu = 0$ by Lemma 2.
So $V_\lambda \neq V_\mu$.

Thus, we have found as many irreps as there are conjugacy classes. Since there can't be more, we must have found all of them. The proof is complete.

Beispiel: Die 2-dimensionale Darstellung von S_3 gehört

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zur Partition $3 = 2+1$, i.e. das Young-Diagramm ist $\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \\ \hline \end{array}$:

Das kanonische Tableau ist $\begin{array}{|c|c|} \hline 1 & 2 \\ \hline 3 & \\ \hline \end{array}$, der zugehörige

Young-Symmetrisierer ist $c_\lambda = (1 + (12))(1 - (13)) = 1 + (12) - (13) - (132)$.

Hierbei bezeichnen runde Klammern Zykeln $(1 \sigma(1) \sigma^2(1) \dots)$, deren

Verknüpfung man z.B. durch die Wirkung auf $N = \{1, 2, 3\}$

bestimmt:

$$(12)(13)\{1, 2, 3\} = (12)\{3, 2, 1\} = \{2, 1, 2\} = (132)\{1, 2, 3\}$$

Benutzen wir für CS_n die geschweifte Notation, so ist

$$Y_\lambda(\{1, 2, 3\}) = \{1, 2, 3\} + \{2, 1, 3\} - \{3, 2, 1\} - \{3, 1, 2\} =: e_1$$

$$Y_\lambda(\{1, 3, 2\}) = (23)\{Y_\lambda(\{1, 2, 3\})\} = \{1, 3, 2\} + \{3, 1, 2\} - \{2, 3, 1\} - \{2, 1, 3\} =: e_2$$

$$(12)e_1 = \{2, 1, 3\} + \{1, 2, 3\} - \{3, 1, 2\} - \{3, 2, 1\} = e_1$$

$$(12)e_2 = \{2, 3, 1\} + \{3, 2, 1\} - \{1, 3, 2\} - \{1, 2, 3\} = -e_2 - e_1$$

$$(23)e_1 = e_2$$







$$(23)e_2 = e_1$$



~~Da $(12), (23) \in S_3$ erzeugen, folgt bereits, dass e_1 und e_2~~

~~$V_{\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \\ \hline \end{array}}$ aufspannen und aus da für $\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \\ \hline \end{array} (23) = \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \\ \hline \end{array} (12) = \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \\ \hline \end{array}$ können~~

~~wird die folgende Identifikation~~

Da (12) und (23) S_3 erzeugen, folgt bereits, dass e_1 und e_2 $V_{\mathbb{R}}$ aufspannen. Aus $\text{tr}(23) = \text{tr}\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = 0$ und $\text{tr}(23(12)) = \text{tr}\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & -1 \\ 0 & -1 \end{pmatrix} = \text{tr}\begin{pmatrix} 0 & -1 \\ 1 & -1 \end{pmatrix} = -1$ folgt die Identifikation mit der 2-dimensionalen Darstellung aus Kapitel 1. Die mit Young-Diagrammen vertierte Charaktertafel ist also jetzt:

			
	1	1	1
	2	0	-1
	1	-1	1

Hätten wir statt dem kanonischen das andere Standardtableau (23)  =  benutzt, so wäre das entsprechende Young-Symmetrisierer

$$c'_{\lambda_{(23)}} = (1 + (13)) (1 - (12)) = 1 + (13) - (12) - (123)$$

und

$$e'_1 = \{1, 2, 3\} + \{3, 2, 1\} - \{2, 1, 3\} - \{2, 3, 1\}$$

$$e'_2 = \{1, 3, 2\} + \{2, 3, 1\} - \{3, 1, 2\} - \{3, 2, 1\} = (23) e_1$$

$$(12) e'_1 = -e'_1 - e'_2$$

$$(12) e'_2 = e'_2$$

	{1,2,3}	{1,3,2}	{2,1,3}	{2,3,1}	{3,1,2}	{3,2,1}
e_1	1	0	1	0	-1	-1
e_2	0	1	-1	-1	1	0
e'_1	1	0	-1	-1	0	1
e'_2	0	1	0	1	-1	-1
$\begin{array}{ c } \hline \square \\ \hline \end{array}$	1	1	1	1	1	1
$\begin{array}{ c } \hline \square \\ \hline \square \\ \hline \end{array}$	1	-1	-1	1	1	-1

Check: e_1, e_2 and e'_1, e'_2 are linearly independent.

* Der Beweis der Aussagen über eine Basis von V_λ und $\text{Hom}(V_\lambda, CS_n) = V_\lambda^*$ an dieser Stelle stellt sich als zu aufwendig heraus.

Die Idee ~~ist~~^{wäre gewesen,} zunächst die lineare Unabhängigkeit des Standardtableaus zu zeigen (was einfacher ist als die Erzeugung) und anschließend kombinatorisch zu verifizieren, dass

$$n! = |S_n| = \sum_{\lambda \text{ Partition von } n} \left(\# \left\{ \begin{array}{l} \text{Standardtableaus} \\ \text{der Form } \lambda \end{array} \right\} \right)^2$$

Siehe z.B. B.Simon, Reps of finite & compact groups
 G.D.James, The representation theory of the symmetric groups.

(aus letzterem Buch stammt wohl auch der folgende Algorithmus für die Charaktertafel, der ohne symmetrischer Funktionen auskommt.)

§3.3. Die Charaktertafel der symmetrischen Gruppe

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Group and representation theory - Lecture 5

Representation theory of the symmetric group II - Characters

In the previous lecture, we reviewed the fact that conjugacy classes $C_{\vec{k}}$ in the symmetric group S_n are uniquely determined by the cycle type

$$\vec{k} = (k_1, k_2, \dots) \quad \sum i k_i = n$$

where $k_i = \#$ cycles of length i .

We then showed how to associate irreducible representations V_{λ} to "Young tableaux of shape λ ", where

$$\lambda = (\lambda_1 \geq \lambda_2 \geq \dots) \quad \sum \lambda_j = n$$

is a partition of n .

$$V_{\lambda} = \mathbb{C} S_n e_{\lambda} = \mathbb{C} S_n p_{\lambda} q_{\lambda}$$

where $p_{\lambda} = \sum_P \pi$ $q_{\lambda} = \sum_Q \sigma$

and P, Q are row, column stabilizer of given Young tableau.

We now wish to calculate the ^{irreducible} characters

$$\chi_{\lambda}(\vec{k}) = \text{tr}_{V_{\lambda}} \sigma, \quad \sigma \in C_{\vec{k}}$$

not quite the same definition, but OK

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As notation and terminology, V_λ is known as "Specht module" (module being synonymous with representation in this course).

The row stabilizes $P = P_\lambda$, whose embedding in S_n depends on tableau T is known as Young subgroup.

- We recall that $|P| = \prod_j \lambda_j!$, and that the size of conjugacy class is given by the formula

$$|C_{\vec{k}}| = \frac{n!}{\text{size of stabilizer}} = \frac{n!}{\prod_i i^{k_i} k_i!}$$

centralizer

- We know from general theory that ~~the~~ irreducible characters are orthogonal w.r.t. particular inner product

$$\sum_{\vec{k}} \bar{\chi}_\lambda(\vec{k}) \frac{|C_{\vec{k}}|}{n!} \chi_\mu(\vec{k}) = \delta_{\lambda\mu}$$

(the formulas imply that χ 's are all real. I'll omit the $\bar{}$)

- We write this as matrix formula. We have already ordered partitions (lexicographically)

$$\lambda > \mu \quad \text{if} \quad \lambda_i - \mu_i > 0 \quad \text{for first non-zero.}$$

Identifying cycle types with partitions allows us to also order \vec{k} 's and even compare χ 's with \vec{k} 's.

With this ordering, we let $X = (X_{\lambda \vec{k}})$ be the character table as a matrix in which rows are indexed by irreps (λ) and columns by conjugacy classes (\vec{k}).

Then, let Σ be diagonal matrix

$$\Sigma_{\vec{k}\vec{l}} = \delta_{\vec{k}\vec{l}} \frac{|C_{\vec{k}}|}{n!}$$

so that orthogonality of characters becomes

$$X \Sigma X^T = 1.$$

The main step towards calculating X is to introduce the following "auxiliary representations" W_λ ~~also~~ also indexed by partitions, but generally reducible. With same notation as before,

$$W_\lambda = \mathbb{C} S_n P_\lambda. \quad \text{i.e. only symmetric the rows.}$$

Calculation of character of W_λ is fairly straightforward. Let $\sigma \in C_{\vec{k}}$. Then

$$\chi_\lambda(\vec{k}) = \text{tr}_{W_\lambda} \sigma = \# \left\{ \tau \in S_n, \sigma \tau = \tau \pi, \text{ some } \pi \in P_\lambda \right\} / \# P_\lambda$$

(in words, we want the number of tableaux such that acting with σ can be repaired by a row permutation, but we identify two such if they are related by a row permutation)

this is,

$$\begin{aligned}
t_{W_\lambda} \sigma &= \# \{ \tau \in S_n, \tau \sigma \tau \in P_\lambda \} / |P_\lambda| \\
&= \frac{|P_\lambda \cap C_{\vec{k}}|}{|P_\lambda|} \cdot |\text{stabilizer of } \sigma|
\end{aligned}$$

$$\psi_\lambda(\vec{k}) = \frac{|P_\lambda \cap C_{\vec{k}}| \cdot n!}{|P_\lambda| |C_{\vec{k}}|}$$

The W_λ is an example of a so-called induced representation (here: the rep of S_n induced by the subgroup P_λ trivial rep of the subgroup P_λ and the formula for $\psi_\lambda(\vec{k})$ is an example of Frobenius reciprocity. See Kapitel 2. See end of lecture 3 which we did not cover in class.

I'll give a more explicit formula for $\psi_\lambda(\vec{k})$ in terms of symmetric functions, below. In fact, most of the literature also does the next step, the decomposition of $\psi_\lambda(\vec{k})$ into irreducibles, in terms of symmetric functions. However this is unnecessary. A noteworthy thing however is that $\psi_\lambda(\vec{k}) = 0$ if $\vec{k} > \lambda$. (in the conventions of before maximal length of cycle in S_{λ_i} is λ_i .)

Therefore, the matrix

(this doesn't play any role actually...)

$$\Psi = (\Psi_\lambda(\vec{k}))$$

is in fact lower triangular.

Now on general grounds we know that

$$W_\lambda \cong \bigoplus_{\mu} V_\mu \oplus K_\lambda^\mu \quad (*)$$

for some non-negative integer degeneracies $K_\lambda^\mu \in \mathbb{N}_{\geq 0}$ known as Kostka numbers.

Claim: The matrix $K = (K_\lambda^\mu)$ is upper triangular and invertible (i.e. $K_\lambda^\mu = 0$ if $\mu < \lambda$ and $K_\lambda^\lambda > 0$.)

The second claim follows immediately from the fact that

$$V_\lambda = W_\lambda q_\lambda.$$

i.e. right-multiplication by q_λ defines a ^{non-zero} map in

$$\text{Hom}_{\mathbb{C}\mathcal{S}_n}(W_\lambda, V_\lambda)$$

To check the first claim, we verify that if $\lambda > \mu$

$$\text{Hom}_{\mathbb{C}\mathcal{S}_n}(W_\lambda, V_\mu) = \text{Hom}_{\mathbb{C}\mathcal{S}_n}(\mathbb{C}\mathcal{S}_n p_\lambda, \mathbb{C}\mathcal{S}_n q_\mu q_\mu) = 0$$

Indeed, to specify such a map, say $f: W_\lambda \rightarrow V_\mu$, it's enough to give $f(p_\lambda) = \sum x_{\mu} p_\mu$, $x \in \mathbb{C}^n$

But since $p_\lambda^2 \propto p_\lambda$ we find

$$\begin{aligned} f(p_\lambda) &= \frac{1}{|P_\lambda|} f(p_\lambda^2) = \frac{1}{|P_\lambda|} P_\lambda f(p_\lambda) \\ &= \frac{1}{|P_\lambda|} P_\lambda \sum x_\mu p_\mu = 0 \quad \text{if } \lambda > \mu \text{ by Lemma 2} \end{aligned}$$

As a consequence of (*),

$$\Psi_\lambda(\vec{k}) = \sum K_\lambda^\mu \chi_\mu(\vec{k})$$

i.e.

$$\Psi = KX$$

Therefore,

$$\Psi \Sigma \Psi^T = K X \Sigma X^T K^T$$

(inner product of characters)
computed in terms of irreducibles)

$$\Psi \Sigma \Psi^T = K K^T$$

It's not hard to check that an upper triangular matrix K with positive diagonal entries is uniquely determined by this equation.
Integrality follows from general principles.

At the end,

$$X = K^{-1} \Psi$$

and we are done.

Addendum: Symmetric functions

• First useful observation is the generating function for the number of partitions:

$$\sum_{n=0}^{\infty} N(n)q^n = \prod_{k=1}^{\infty} \sum_{n=0}^{\infty} q^{nk}$$

$$= \prod_{k=1}^{\infty} \frac{1}{1-q^k}$$

- closely related to Dedekind η -function.

$$N(n) \sim \frac{1}{4\sqrt{3}n} e^{\frac{\pi\sqrt{213}}{3}n}$$

• Second, let's calculate $|P_{\lambda} \cap C_k^{\rightarrow}|$. We consider conjugacy classes in

$$P_{\lambda} = \prod_{j=1}^r S_{\lambda_j}$$

they are labelled by collections (r_{ij}) s.t.

$$\lambda_j = \sum_i i r_{ij}$$

To get an element in conjugacy class C_k^{\rightarrow} of S_n , we need

$$\sum_j r_{ij} = k_i$$

Then

$$|P_{\lambda} C_{\vec{k}}| = \sum_{(r)} \prod_j \frac{\lambda_j!}{\prod_i i^{r_{ij}} r_{ij}!}$$

So

$$\psi_{\lambda}(\vec{k}) = \frac{1}{\prod \lambda_j!} \prod_i i^{k_i} k_i! \sum_{(r)} \prod_j \frac{\lambda_j!}{\prod_i i^{r_{ij}} r_{ij}!}$$

$$= \sum_{(r)} \prod_i \frac{k_i!}{\prod_j i^{r_{ij}} r_{ij}!} \quad \begin{matrix} \sum_j r_{ij} = k_i \\ \sum_i i r_{ij} = \lambda_j \end{matrix}$$

Frobenius observation is that this combination of multinomial coefficients is also the coefficient of

$$x^{\lambda} = x_1^{\lambda_1} x_2^{\lambda_2} \dots x_s^{\lambda_s}$$

in $\underbrace{(x_1 + \dots + x_s)^{k_1} (x_1^2 + \dots + x_s^2)^{k_2} \dots}$

$$\psi_{\lambda}(\vec{k}) = [P^{(\vec{k})}]_{\lambda} \quad \text{— coefficient of } x^{\lambda}$$

which is the beginning of a long story involving symmetric functions, at the end of which one proves the Frobenius character formula

$$\chi_\lambda(\vec{k}) = [\Delta P^{(\vec{k})}]_l$$

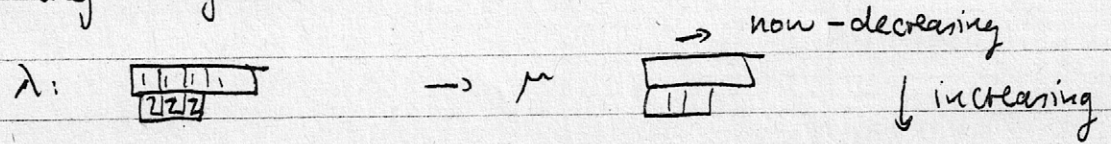
$$= (\text{coefficient of } x^l \text{ in } \Delta \cdot P^{(\vec{k})})$$

where $P^{(\vec{k})}$ as above,

$$\Delta = \prod_{i < j} (x_i - x_j) \quad (\text{Vandermonde determinant of } \begin{pmatrix} 1 & x_1 & x_1^2 & \dots & x_1^{s-1} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & x_2 & \vdots & \vdots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & x_s & \dots & \dots & x_s^{s-1} \end{pmatrix})$$

$$\text{and } l = (\lambda_1 + s - 1, \lambda_2 + s - 2, \dots, \lambda_{s-1} + 1, \lambda_s)$$

The gist of it is to define Kostka numbers combinatorially as ways of filling up Young diagrams



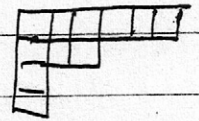
and then show that corresponding decomposition of $\psi_\lambda(\vec{k})$ is irreducible characters.

(what I described is obviously false, though it might give less information.)

That $K_\lambda^\mu = \text{Hom}(V_\mu, W_\lambda)$ follows as corollary.)

Among the consequences of the Frobenius character formula is the so-called Hook length formula.

$$\dim V_\lambda = \frac{n!}{\prod_{\text{boxes}} (\text{Hook length})}$$



7	4	2	1
4	1		
2			
1			

$$\frac{8!}{7 \cdot 4 \cdot 2 \cdot 4 \cdot 2} = 6 \cdot 5 \cdot 3 = 90.$$

(whose equality with number of standard tableaux is hard to prove, combinatorially.)

as well as a variety of other combinatorial results (Murnaghan-Nakayama rule, etc.)